



# Ruijsenaars duality for $B, C, D$ Toda chains

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# GENERALIZED TODA CHAINS

- $H(p, q) = \sum_{i=1}^n \frac{p_i^2}{2} + \sum_{i=1}^{n-1} e^{2(q_i - q_{i+1})}$  — A-type Toda chain
  - other root systems
- $$H(p, q) = \frac{1}{2}(p, p) + \sum_{\alpha \in \Pi} c_\alpha e^{2\alpha(q)} \quad (p, q) \in \mathfrak{h}, \alpha \in \Pi \subset \mathfrak{h}^*$$
- simple roots

- classical integrable system
- n independent Poisson-commuting functions  $\{H_i, H_j\} = 0$

$$H_1(p, q) = \sum_i p_i$$

$$H_2(p, q) = H(p, q)$$

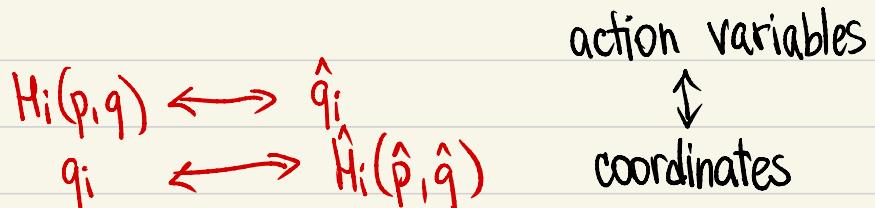
$$H_l(p, q) = \frac{1}{l} \sum_i p_i^l + (\text{lower terms in } p)$$

# RUIJSENAARS DUALITY

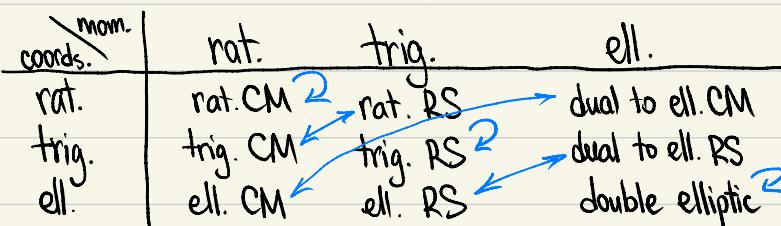
- classical integrable system:  $\{H_i, H_j\} = 0 \quad i, j = 1, \dots, n$
- another Poisson-commuting set:  $\{q_i, q_j\} = 0 \quad i, j = 1, \dots, n$

Ruijsenaars duality:

two integrable systems  
on the same phase space



Dualities square for  
Calogero - Ruijsenaars  
integrable systems



# GOLDFISH MODELS

- Ruijsenaars dual to Toda chains

$$\hat{H}(\hat{p}, \hat{q}) = \sum_{i=1}^n e^{2\hat{p}_i} \prod_{j \neq i} \frac{1}{\hat{q}_i - \hat{q}_j}$$

momentum in exp-trigonometric coordinates – rational dependence (opposite to the Toda case)

trigonometric CM system

$$H(p, q) = \sum_{i=1}^n \frac{p_i^2}{2} + \sum_{i \neq j} \frac{g^2}{2 \sin^2(q_j - q_i)}$$

Ruijsenaars duality

rational RS system

$$\hat{H}(\hat{p}, \hat{q}) = \sum_{i=1}^n e^{2\hat{p}_i} \prod_{j \neq i} \frac{\hat{q}_i - \hat{q}_j + \eta}{\hat{q}_i - \hat{q}_j}$$

Inozemtsev limit

$$H(p, q) = \sum_{i=1}^n \frac{p_i^2}{2} + \sum_{i=1}^{n-1} e^{2(q_i - q_{i+1})}$$

Strong coupling limit

(open) Toda chain

$$\hat{H}(\hat{p}, \hat{q}) = \sum_{i=1}^n e^{2\hat{p}_i} \prod_{j \neq i} \frac{1}{\hat{q}_i - \hat{q}_j}$$

(rational) Goldfish model

# HAMILTONIAN REDUCTION

- $(M, \omega)$  — symplectic manifold with Hamiltonian action of Lie group  $G$   
 $\text{Ad}_G^*$ -equivariant momentum map  $\mu: M \rightarrow \mathfrak{g}^*$   $\langle \mu(x), \xi \rangle = H_\xi(x)$

- reduced phase space  $M_\lambda = \mu^{-1}(\lambda) / \text{Stab}^{\text{Ad}^*}(\lambda)$  for fixed  $\lambda \in \mathfrak{g}^*$

- symplectic form  $\omega_\lambda: \mu^{-1}(\lambda) \xrightarrow{i} M$   $i^* \omega = \pi^* \omega_\lambda$

- source of integrable systems:

simple functions  $f, g \in C(M)^G$   
large phase space  $\{f, g\}_M = 0$



nontrivial functions  
on reduced space  $\tilde{f}, \tilde{g} \in C(M_\lambda)$   
 $\{\tilde{f}, \tilde{g}\}_{M_\lambda} = 0$

# LIE GROUPS NOTATIONS

- $G$  - matrix complex Lie group  $\mathfrak{g} = \mathfrak{n}_+ \oplus \mathfrak{h} \oplus \mathfrak{n}_-$
  - $N_{\pm}$  - unipotent Lie groups  $\text{Lie}(N_{\pm}) = \mathfrak{n}_{\pm}$
  - $K$  - maximal compact Subgroup  $\text{Lie}(K) = \mathfrak{k}$
  - type A  $G = \text{GL}(n, \mathbb{C})$
  - type C  $G = \text{Sp}(2n, \mathbb{C}) = \{g \in \text{GL}(2n, \mathbb{C}) \mid g \Omega g^T = \Omega\}$
- $\Omega = \begin{pmatrix} 0 & P \\ -P & 0 \end{pmatrix}_{2n \times 2n}$  skew-symmetric form
- $P = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}_{n \times n}$
- $N_+ = \left\{ \begin{pmatrix} 1 & * \\ 0 & 1 \end{pmatrix} \right\}$  upper-triangular  
1's on the diagonal
- $N_+ = \left\{ \begin{pmatrix} A & B \\ 0 & P(A^T)^{-1}P^{-1} \end{pmatrix} \mid A = \begin{pmatrix} 1 & * \\ 0 & 1 \end{pmatrix} \right\}$
- $K = U(n)$  ← both unitary →  $K = U(2n) \cap \text{Sp}(2n, \mathbb{C})$

# COTANGENT BUNDLE

- left translations trivialization  $T^*G = G \times \mathfrak{g}^* \ni (g, X)$

- canonical symplectic structure  $\omega = \text{Tr}(X g^{-1} dg) \leftarrow \{X_a, X_b\} = -f_{ab}^c X_c \leftarrow$   
 $\omega = dd^* \leftarrow \{g, g\} = 0 \leftarrow g$  Lie-algebraic  
matrix elements  
commute

- left and right Hamiltonian actions

$$L_h: (g, X) \mapsto (hg, X)$$

$$\mu_L(g, X) = gXg^{-1}$$

$$R_h: (g, X) \mapsto (gh^{-1}, hXh^{-1})$$

$$\mu_R(g, X) = -X$$

- subgroups actions

$N_+$  left action

$$\mu_L(g, X) = \text{Pr}_{N_+^*}(gXg^{-1}) \in \mathfrak{n}_+^* \simeq \mathfrak{n}_-$$

$K$  right action

$$\mu_R(g, X) = -\text{Pr}_{K^*}(X) \in \mathfrak{k}^* \simeq \mathfrak{k}$$

use  $\text{Tr}$  to identify

# REDUCED PHASE SPACE

- momentum maps

$$M_L(g, X) = \text{Pr}_{N_-}(gXg^{-1}) = \lambda_L = \sum_{\alpha \in \Delta^+} e_{-\alpha}$$

↑  
simple roots

$$M_R(g, X) = -\text{Pr}_{N_+}(X) = \lambda_R = 0$$

maximal possible stabilizer —  $N_+ \times K$

$$n \left( \sum_{\alpha \in \Delta^+} e_{-\alpha} \right) h^{-1} = \sum_{\alpha \in \Delta^+} e_{-\alpha} \quad \forall n \in N_+$$

- reduced phase space  $M_\lambda = \{(g, X) \mid M_L = \lambda_L, M_R = \lambda_R\} / N_+ \times K$   $\dim_R M_L = 2 \text{rk } G$
- invariant commuting functions

$$I_\ell = \text{Tr } X^\ell \quad \xrightarrow{\text{rk } G \text{ independent}}$$

$$\{\text{Tr } X^\ell, \text{Tr } X^m\} = 0$$

as Casimirs  $\{X_a, X_b\} = -f_{ab}^c X_c$

$$J_\ell = m_\ell(g g^+)$$

lowest-right  $\ell \times \ell$  minors

$$\{g, g\} = 0 \Rightarrow \{J_\ell, J_m\} = 0$$

$$gg^+ \mapsto n g k^{-1} \cdot \underbrace{(k^{-1})^+}_{=1, \text{k-unitary}} g^+ n^+ = n \cdot gg^+ \cdot n^+ \Rightarrow \text{invariant}$$

upper-tr.  
with 1's  
diagonal

lower-tr.  
with 1's  
diagonal

# TODA GAUGE

reduced phase space  $\{Pr_{\beta}(X) = 0, Pr_n(gXg^{-1}) = \sum_{\alpha \in \Pi} e_{\alpha}\}/N_+ \times K$

- Toda gauge — use  $N_+ \times K$  to diagonalize  $g \mapsto ngk^{-1}$   
Iwasawa decomposition  $G \rightarrow N_+ \times A \times K$ ,  $A$  — diagonal matrices  

$$g = \exp\left(\sum_{i=1}^{\text{rk } G} q_i h_i\right) \quad q_i \in \mathbb{R}$$
- Solve momentum equations

$$Pr_{\beta}(X) = 0 \Rightarrow X = p + \sum_{\alpha > 0} X_{\alpha} (e_{\alpha} + e_{-\alpha}), \quad Pr_n(gXg^{-1}) = 0 \Rightarrow \begin{cases} X_{\alpha} = e^{\alpha(q)}, \alpha \in \Pi \\ X_{\alpha} = 0, \alpha \notin \Pi \end{cases}$$

- Toda chain Lax operator  $X = p + \sum_{\alpha \in \Pi} e^{\alpha(q)} (e_{\alpha} + e_{-\alpha})$

$I_l = \text{Tr } X^l$  — Toda chain Hamiltonians

$J_l = m_l(gg^+)$  — simple functions of Toda coordinates

# MOSER GAUGE

reduced phase space  $\{ \text{Pr}_\beta(X) = 0, \text{Pr}_{n_-}(gXg^{-1}) = \sum_{\alpha \in \Pi} e_{-\alpha} \} / N \times K$

- Moser gauge — use  $X \mapsto kXk^{-1}$  to diagonalize  $X = \sum \hat{q}_i h_i$  ← always possible  $\text{Pr}_\beta(X) = 0$
- $N_+$  gauge freedom — make  $g$  lower-triangular using Gauss decomposition  $g \in B_-$  ← opposite Borel subgroup
- Solve momentum equations

$$\text{Pr}_{n_-}(gXg^{-1}) = \sum_{\alpha \in \Pi} e_{-\alpha} \Rightarrow gXg^{-1} - X = \sum_{\alpha \in \Pi} e_{-\alpha} \Rightarrow gX = Xg + \sum_{\alpha \in \Pi} e_{-\alpha} \cdot g$$

$\xrightarrow[g \in B_- \quad X \in \mathfrak{h}]{} gXg^{-1} \in \mathfrak{h}_- = \mathfrak{h} \oplus \mathfrak{n}_-, \quad \text{Pr}_{\mathfrak{h}}(gXg^{-1}) = X$

- $I_\ell = \text{Tr } X^\ell$  — simple functions of dual coordinates  $J_\ell = m_\ell(gg^t)$  — dual Hamiltonians

# TYPE A : SOLVING MOMENTUM EQ

Moser gauge  $X = \sum_{i=1}^{nG} \hat{q}_i h_i$  - diagonal matrix, real-valued  $\hat{q}_i \in \mathbb{R}$

for  $G = GL(n, \mathbb{C})$   $X = \begin{pmatrix} \hat{q}_1 & & 0 \\ & \hat{q}_2 & \\ 0 & \ddots & \hat{q}_n \end{pmatrix}$   $\lambda_L = \sum_{\alpha \in \Gamma} e_{-\alpha} = \begin{pmatrix} 0 & & \\ \vdots & 0 & \\ 0 & \ddots & 0 \end{pmatrix}$

- momentum equations:  $gX - Xg = \lambda_L g \Rightarrow g_{ij}(\hat{q}_j - \hat{q}_i) = g_{i-1,j}$  ( $i \neq j$ )

*no constraints on  
diagonal elements of g*

*recursive formula for g matrix elements  
Starting from diagonal elements*

- denote  $g_{ii} = \hat{q}_i$ , solution of momentum equations

$$g_{ij} = \hat{q}_j \cdot \prod_{k=j+1}^i \frac{1}{\hat{q}_j - \hat{q}_k} \quad (i > j)$$

*lower-triangular*

# TYPE C : SOLVING MOMENTUM EQ

Moser gauge  $X = \sum_{i=1}^{n+6} \hat{q}_i h_i$  - diagonal matrix, real-valued  $\hat{q}_i \in \mathbb{R}$

for  $G = Sp(2n, \mathbb{C})$   $X = \begin{pmatrix} \hat{q}_1 & & & & & \\ & \ddots & \hat{q}_n & & & \\ & & -\hat{q}_n & & & \\ & & & \ddots & & \\ 0 & & & & \ddots & -\hat{q}_1 \\ & & & & & \end{pmatrix}$   $\lambda_L = \begin{pmatrix} 0 & & & & & \\ & 1 & & & & \\ & & -1 & 0 & & \\ & & & 1 & 0 & \\ 0 & & & & -1 & \\ & & & & & i0 \end{pmatrix}$

- fix diagonal elements  $g = \begin{pmatrix} A & 0 \\ C & P(A^T)^{-1}P \end{pmatrix}, A = \begin{pmatrix} \hat{a}_1 & 0 \\ * & \ddots & \hat{a}_n \end{pmatrix}$

- recursive relations from momentum equation

$$\begin{cases} (\hat{q}_j - \hat{q}_i) A_{ij} = A_{i-1,j} \\ (\hat{q}_{n+i-j} + \hat{q}_j) C_{ij} - \delta_{i,1} A_{nj} + C_{i-1,j} = 0 \end{cases}$$

- Solution  $A_{ij} = \hat{a}_j \prod_{k=j+1}^i \frac{1}{\hat{q}_j - \hat{q}_k}, C_{ij} = (-1)^{i-1} \hat{a}_j \prod_{k=j+1}^n \frac{1}{\hat{q}_{j-k}} \prod_{m=1}^i \frac{1}{\hat{q}_{n+1-m} + \hat{q}_j}$

# TYPE A : DUAL HAMILTONIANS

- Hamiltonians using Cauchy-Binet formula

$$J_\ell = \sum_{|S|=\ell} \hat{a}_{S_i}^2 \times \prod_{r=1}^{\ell} \left( \prod_{\substack{j=s_r+1 \\ j \notin S}}^n \frac{1}{\hat{q}_{s_r} - \hat{q}_j} \right)^2$$

- Change of variables

$$e^{2\hat{p}_i} = \hat{a}_i^2 \cdot \prod_{j=i+1}^n (\hat{q}_i - \hat{q}_j) \cdot \prod_{j=1}^{i-1} \frac{1}{\hat{q}_i - \hat{q}_j}$$

for more symmetric form of Hamiltonians

$$J_\ell = \sum_{|S|=\ell} \prod_{i=1}^{\ell} e^{2\hat{p}_{S_i}} \cdot \prod_{\substack{a \in S \\ b \notin S}} \frac{1}{\hat{q}_a - \hat{q}_b}$$

$$\text{e.g. } J_1 = \sum_{i=1}^n e^{2\hat{p}_i} \prod_{j \neq i} \frac{1}{\hat{q}_i - \hat{q}_j}$$

- strong coupling limit of rational Ruijsenaars-Schneider

$$H_\ell = \sum_{|S|=\ell} \prod_{i=1}^{\ell} e^{2\hat{p}_{S_i}} \cdot \prod_{\substack{a \in S \\ b \notin S}} \frac{\hat{q}_a - \hat{q}_b + \eta}{\hat{q}_a - \hat{q}_b}$$

# TYPE C: DUAL HAMILTONIANS

- $g = \begin{pmatrix} A & 0 \\ C & P(A^T)^{-1}P \end{pmatrix}, \quad A = \begin{pmatrix} \hat{a}_1 & 0 \\ * & \ddots & \hat{a}_n \end{pmatrix}$
- only these matrix elements appear in  $J_\ell = m_\ell(g g^+), \quad 1 \leq \ell \leq n$
- $(P(A^T)^{-1}P)_{ij} = (-1)^{i-j} \hat{a}_{n+1-j}^{-1} \prod_{k=j+1}^i \frac{1}{\hat{q}_{n+1-k} - \hat{q}_{n+1-j}}$
- $C_{ij} = (-1)^{j-i} \hat{a}_j \prod_{k=j+1}^n \frac{1}{\hat{q}_j - \hat{q}_k} \prod_{m=1}^i \frac{1}{\hat{q}_{n+1-m} + \hat{q}_j}$
- Cauchy-Binet + A type calculations  $\rightarrow$  Hamiltonians
- $J_1$  after change of variables  $J_1 = \sum_{i=1}^n (e^{2\hat{p}_i} + e^{-2\hat{p}_i}) \frac{1}{2\hat{q}_i} \prod_{j \neq i} \frac{1}{(\hat{q}_i - \hat{q}_j)(\hat{q}_i + \hat{q}_j)}$

# REVIEW OF ANSWERS

Toda system

$$H = \sum_{i=1}^n \frac{p_i^2}{2} + \sum_{i=1}^{n-1} e^{2(q_i - q_{i+1})}$$

$$H = \sum_{i=1}^n \frac{p_i^2}{2} + \sum_{i=1}^{n-1} e^{2(q_i - q_{i+1})} + e^{2q_n}$$

$$H = \sum_{i=1}^n \frac{p_i^2}{2} + \sum_{i=1}^{n-1} e^{2(q_i - q_{i+1})} + \frac{1}{2} e^{4q_n}$$

$$H = \sum_{i=1}^n \frac{p_i^2}{2} + \sum_{i=1}^{n-1} e^{2(q_i - q_{i+1})} + e^{2(q_{n-i} + q_n)}$$

A

Goldfish model

$$\hat{H} = \sum_{i=1}^n e^{2\hat{p}_i} \prod_{j \neq i} \frac{1}{\hat{q}_i - \hat{q}_j}$$

B

$$\hat{H} = \sum_{i=1}^n (e^{2\hat{p}_i} + e^{-2\hat{p}_i}) \frac{1}{2\hat{q}_i^2} \prod_{j \neq i} \frac{1}{(\hat{q}_i - \hat{q}_j)(\hat{q}_i + \hat{q}_j)} + \prod_{j=1}^n \frac{1}{\hat{q}_j^2}$$

C

$$\hat{H} = \sum_{i=1}^n (e^{2\hat{p}_i} + e^{-2\hat{p}_i}) \frac{1}{2\hat{q}_i} \prod_{j \neq i} \frac{1}{(\hat{q}_i - \hat{q}_j)(\hat{q}_i + \hat{q}_j)}$$

D

$$\hat{H} = \sum_{i=1}^n (e^{2\hat{p}_i} + e^{-2\hat{p}_i}) \prod_{j \neq i} \frac{1}{(\hat{q}_i - \hat{q}_j)(\hat{q}_i + \hat{q}_j)}$$