Affine Weyl groups and non-abelian discrete systems

Irina Bobrova

MPI for Mathematics in the Sciences, Leipzig, Germany

July 10, 2024

"RTISART - 2024", Beijing

based on arXiv:2403.18463 (IB)

A short overview of the topic

Main examples: Painlevé equations.

- Problem: define new functions by an ODE of the *m*th order with properties that generalize those of elliptic functions. [L. Fuchs], [H. Poincaré]
- Painlevé property: the general solution of an ODE has no critical movable points.
- m = 2: six classes defining the Painlevé transcendents. [Painlevé, 1902], [Gambier, 1910]
- Painlevé equations being one of the important objects in mathematics and mathematical physics have various non-commutative (NC) analogs: quantum [Nagoya, 2004], matrix differential [Kawakami, 2015], NC "differential" [Bobrova and Sokolov, 2023b], matrix difference [Cassatella-Contra et al., 2014], NC difference [Bobrova, Retakh, Rubtsov, Sharygin, 2024]
- Some of them are connected with integrable non-abelian PDEs [Olver and Sokolov, 1998] and P∆Es [Adler, 2020], Riemann-Hilbert problem [Cafasso and Manuel, 2014], orthogonal polynomials [Cafasso et al., 2018], Calogero systems [Bertola et al., 2018], and etc.
- In the commutative case, discrete Painlevé equations have been studied in a series of papers by B. Gramaticos and A. Ramani since 1990s, but without understanding the whole picture.
- The latter was clarified by H. Sakai in his famous paper [Sakai, 2001], whose geometric method was inspired by a series of K. Okamoto's papers.
- NC birational geometry and Painlevé equations. [Okounkov and Rains, 2015], [Rains, 2019]
- An application of the affine Weyl groups to discrete Painlevé equations. [Noumi and Yamada, 1998], [Noumi and Yamada, 2000]
- Symmetries of the Hamiltonian matrix Painlevé equations. [Bershtein et al., 2023]

A short overview of the topic

Main examples: Painlevé equations.

- Problem: define new functions by an ODE of the *m*th order with properties that generalize those of elliptic functions. [L. Fuchs], [H. Poincaré]
- Painlevé property: the general solution of an ODE has no critical movable points.
- m = 2: six classes defining the Painlevé transcendents. [Painlevé, 1902], [Gambier, 1910]
- Painlevé equations being one of the important objects in mathematics and mathematical physics have various non-commutative (NC) analogs: quantum [Nagoya, 2004], matrix differential [Kawakami, 2015], NC "differential" [Bobrova and Sokolov, 2023b], matrix difference [Cassatella-Contra et al., 2014], NC difference [Bobrova, Retakh, Rubtsov, Sharygin, 2024]
- Some of them are connected with integrable non-abelian PDEs [Olver and Sokolov, 1998] and P∆Es [Adler, 2020], Riemann-Hilbert problem [Cafasso and Manuel, 2014], orthogonal polynomials [Cafasso et al., 2018], Calogero systems [Bertola et al., 2018], and etc.
- In the commutative case, discrete Painlevé equations have been studied in a series of papers by B. Gramaticos and A. Ramani since 1990s, but without understanding the whole picture.
- The latter was clarified by H. Sakai in his famous paper [Sakai, 2001], whose geometric method was inspired by a series of K. Okamoto's papers.
- ▶ NC birational geometry and Painlevé equations. [Okounkov and Rains, 2015], [Rains, 2019]
- An application of the affine Weyl groups to discrete Painlevé equations. [Noumi and Yamada, 1998], [Noumi and Yamada, 2000]
- Symmetries of the Hamiltonian matrix Painlevé equations. [Bershtein et al., 2023]

A short overview of the topic

Main examples: Painlevé equations.

- Problem: define new functions by an ODE of the *m*th order with properties that generalize those of elliptic functions. [L. Fuchs], [H. Poincaré]
- Painlevé property: the general solution of an ODE has no critical movable points.
- m = 2: six classes defining the Painlevé transcendents. [Painlevé, 1902], [Gambier, 1910]
- Painlevé equations being one of the important objects in mathematics and mathematical physics have various non-commutative (NC) analogs: quantum [Nagoya, 2004], matrix differential [Kawakami, 2015], NC "differential" [Bobrova and Sokolov, 2023b], matrix difference [Cassatella-Contra et al., 2014], NC difference [Bobrova, Retakh, Rubtsov, Sharygin, 2024]
- Some of them are connected with integrable non-abelian PDEs [Olver and Sokolov, 1998] and P∆Es [Adler, 2020], Riemann-Hilbert problem [Cafasso and Manuel, 2014], orthogonal polynomials [Cafasso et al., 2018], Calogero systems [Bertola et al., 2018], and etc.
- In the commutative case, discrete Painlevé equations have been studied in a series of papers by B. Gramaticos and A. Ramani since 1990s, but without understanding the whole picture.
- The latter was clarified by H. Sakai in his famous paper [Sakai, 2001], whose geometric method was inspired by a series of K. Okamoto's papers.
- ▶ NC birational geometry and Painlevé equations. [Okounkov and Rains, 2015], [Rains, 2019]
- An application of the affine Weyl groups to discrete Painlevé equations. [Noumi and Yamada, 1998], [Noumi and Yamada, 2000]
- Symmetries of the Hamiltonian matrix Painlevé equations. [Bershtein et al., 2023]

Our goals & results

Global goal

Generalize Sakai's approach to the non-commutative case in order to obtain a full classification of NC versions for the discrete Painlevé equations.

Local goal

 Generalize Noumi-Yamada's approach to the non-commutative case in order to obtain NC versions of the *d*-Painlevé equations.

Results

- Extended representations of the affine Weyl groups W and NC discrete systems.
- NC analogs for the d-Painlevé equations obtained by using extended representations of W. Note that the commutative degeneration scheme holds in the NC case.
- ► NC dressing chains in the Noumi-Yamada variables:
 - ► Lax pairs,
 - Bäcklund transformations,
 - a description of the related discrete dynamics.

A NC version of the $d-P(E_7)$

NC setting

- Consider an associative unital division ring $\mathcal R$ over the field $\mathbb C$ equipped with a derivation.
- We assume that all greek letters belong to the field \mathbb{C} , while the elements f_i are from \mathcal{R} . We will often call f_i as *functions*.
- ▶ The derivation $d_t : \mathcal{R} \to \mathcal{R}$ of the ring \mathcal{R} is a \mathbb{C} -linear map satisfying the Leibniz rule. We also assume that there is a central element t such that $d_t(t) = 1$ and for any $\alpha \in \mathbb{C}$ we have $d_t(\alpha) = 0$. Here and below we identify the unit of the field with the unit of the ring.
- For the brevity we denote $d_t(f_i) = \dot{f}_i$, $d_t^2(f_i) = \ddot{f}_i$, and so on.
- Note that we have an involution on R called the transposition τ, which acts trivially on the generators of R and for any elements F, G ∈ R we have τ(F G) = τ(G) τ(F). This involution can be naturally extended to the matrices over R.

Remark. We would rather not specify the generators of the ring \mathcal{R} in order to avoid an overloaded description of a skew field (see [Cohn, 1995]). Instead, we encourage to think of the ring \mathcal{R} as a generalization of rational functions over \mathbb{C} to a non-commutative case.

The P_2 equation (1)

► Consider the P₂ system [Retakh and Rubtsov, 2010] (see also [Adler and Sokolov, 2021])

$$\begin{cases} \dot{q} &= -q^2 + p - \frac{1}{2}t, \\ \dot{p} &= qp + pq + \alpha_1. \end{cases}$$
 P₂

• Here we can assume that t is also an element of \Re such that $\dot{t} = 1$.

• Let
$$\alpha_0 + \alpha_1 = 1$$
 and $f := -p + 2q^2 + t$.

Its Bäcklund transformations are given below (cf. with [Bershtein et al., 2023])

	α_0	α_1	q	p	t
<i>s</i> 0	$-\alpha_0$	$\alpha_{1}+2\alpha_{0}$	$q - \alpha_0 f^{-1}$	$p - 2\alpha_0 q f^{-1} - 2\alpha_0 f^{-1} q + 2\alpha_0 f^{-2}$	t
<i>s</i> 1	$\alpha_0 + 2\alpha_1$	$-\alpha_1$	$q + \alpha_1 p^{-1}$	p	t
π	α ₁	α_0	-q	$-p+2q^2+t$	t

• These elements form an extended affine Weyl group of type $A_1^{(1)}$:

$$\widetilde{W}(A_1^{(1)}) = \langle s_0, s_1; \pi \rangle,$$

$$s_i^2 = 1, \qquad \pi^2 = 1, \qquad \pi s_i = s_{i+1}\pi, \qquad i \in \mathbb{Z}/_{2\mathbb{Z}}.$$
(1)

The P_2 equation (2)

	α0	α_1	q	p	t
<i>s</i> 0	$-\alpha_0$	$\alpha_{1}+2\alpha_{0}$	$q - \alpha_0 f^{-1}$	$p - 2\alpha_0 q f^{-1} - 2\alpha_0 f^{-1} q + 2\alpha_0 f^{-2}$	t
s_1	$\alpha_0 + 2\alpha_1$	$-\alpha_1$	$q + \alpha_1 p^{-1}$	p	t
π	α_1	α_0	- <i>q</i>	$-p+2q^2+t$	t

Consider the translation operator T = s₁π. It acts on the parameters according to the formula below and forms a lattice on a line:

$$\Gamma(\alpha_0, \alpha_1) = (\alpha_0 - 1, \alpha_1 + 1).$$
(2)

The q and p variables change as follows

 $\bar{q} = s_1 \pi(q) = -s_1(q) = -q - \alpha_1 p^{-1}, \quad \bar{p} = s_1 \pi(p) = s_1(-p + 2q^2 + t) = -p + 2\bar{q}^2 + t.$

• So, we obtain the system for $T(q) = \bar{q}$, $T(p) = \bar{p}$

$\bar{\alpha}_{0} = \alpha_{0} - 1,$	$\bar{\alpha}_1 = \alpha_1 + 1,$	d-P(<i>E</i> ₇)
$\bar{q} + q = -\alpha_1 p^{-1},$	$ar{p}+p=t+2ar{q}^2.$	

It generalizes to the non-commutative case the d-P(E₇) equation from [Sakai, 2001] (p. 206).
 It reduces to the following second-order difference equation for Tⁿ(q) = q_n:

$$\alpha_{1,n} (q_{n+1} + q_n)^{-1} + \alpha_{1,n-1} (q_n + q_{n-1})^{-1} = -2q_n^2 - t, \quad \alpha_{1,n} = \alpha_1 + n. \quad \text{alt-d-P}_1$$

The d-P(E_7) system: a Lax pair

$$\begin{split} \bar{\alpha}_0 &= \alpha_0 - 1, \qquad \bar{\alpha}_1 = \alpha_1 + 1, \\ \bar{q} + q &= -\alpha_1 p^{-1}, \qquad \bar{p} + p = t + 2\bar{q}^2. \end{split} \quad \mathsf{d-P}(E_7)$$

One may also consider the corresponding non-commutative discrete linear problem

$$\begin{cases} \partial_{\lambda} Y &= \mathcal{A} Y, \\ \bar{Y} &= \mathcal{B} Y. \end{cases}$$
(3)

• A Lax pair for the $d-P(E_7)$ is given by

$$\begin{aligned} \mathcal{A} &= \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \lambda^2 + \begin{pmatrix} 0 & 1 \\ 2p & 0 \end{pmatrix} \lambda + \begin{pmatrix} -p + \frac{1}{2}t & -q \\ 2pq + 2\alpha_1 & p - \frac{1}{2}t \end{pmatrix}, \\ \mathcal{B} &= \begin{pmatrix} -2 & 0 \\ 0 & 0 \end{pmatrix} \lambda + \begin{pmatrix} -2q & -1 \\ -2\bar{p} & 0 \end{pmatrix}, \end{aligned}$$

where $t, \lambda \in \mathcal{Z}(\mathcal{R})$.

ų

Note that the compatibility condition is satisfied, since the commutator [p, q] is invariant under the map

$$\psi: \mathfrak{R}^2 \to \mathfrak{R}^2, \quad (q,p) \mapsto (\bar{q},\bar{p}) = \left(-p + t + 2q^2, -q - \bar{\alpha}_1(-p + t + 2q^2)^{-1}\right).$$
 (4)

• Once $t \in \mathbb{R}$, the commutator [p, q] is no longer a conserved quantity.

The d-P(E_7) system: a continuous limit

$$\bar{\alpha}_0 = \alpha_0 - 1, \qquad \bar{\alpha}_1 = \alpha_1 + 1,$$

$$\bar{q} + q = -\alpha_1 p^{-1}, \qquad \bar{p} + p = t + 2\bar{q}^2.$$

$$d-P(E_7)$$

Remark. One may consider a non-commutative analog for the continuous limit as follows. For the simplicity, suppose that we have a difference equation for the functions f_n . One can take the change of variables with the commutative parameter ε

$$z = \varepsilon n$$
 (5)

supplemented by the maps

$$f_n = F, \qquad f_{n+k} = F + k \varepsilon \dot{F} + \frac{1}{2} k^2 \varepsilon^2 \ddot{F} + O(\varepsilon^3).$$
(6)

The latter must be chosen in such a way that the limit $\varepsilon \to 0$ exists.

By using the formulas

$$\begin{aligned} q &= 1 + \varepsilon^2 \, Q - \frac{1}{6} \, \varepsilon^3 \, P, \quad p = -2 + 2\varepsilon^2 \, Q + \frac{2}{3} \, \varepsilon^3 \, P, \quad t = -6 + \frac{1}{3} \, \varepsilon^4 \, T, \quad \alpha_1 = 4 + \frac{2}{3} \, \varepsilon^4 \, T, \\ \text{the d-P}(E_7) \text{ has the P}_1 \text{ system in the limit } \varepsilon \to 0: \end{aligned}$$

$$\begin{cases} \dot{q} &= p, \\ \dot{p} &= 6q^2 + t. \end{cases}$$
 P₁

The $d-P(E_7)$ system: a Hamiltonian form

$$\begin{split} \bar{\alpha}_0 &= \alpha_0 - 1, \qquad \bar{\alpha}_1 = \alpha_1 + 1, \\ \bar{q} + q &= -\alpha_1 p^{-1}, \qquad \bar{p} + p = t + 2\bar{q}^2. \end{split} \quad \mathsf{d}\mathsf{-P}(\mathcal{E}_7)$$

Remark. Here we use non-commutative partial derivatives introduced in [Kontsevich, 1993]. Let $f = f(q, p) \in \mathbb{R}$. Its non-commutative derivatives $\partial_q f$, $\partial_p f$ are defined by the identity

$$df = \partial_q f \, dq + \partial_p f \, dp, \tag{7}$$

where it is assumed that additional non-commutative symbols dq, dp are moved to the right by cyclic permutations of generators in monomials.

Similar to [Veselov, 1991], we call a difference discrete system Hamiltonian if there exists an element $H = H(q, \bar{p}) \in \mathcal{R}$ such that the system can be rewritten in the form

$$p = \partial_q H, \qquad \bar{q} = \partial_{\bar{p}} H.$$
 (8)

- For the d-P(E_7) system, a Hamiltonian is $H = -q \bar{p} + tq + \frac{2}{3}q^3 \bar{\alpha}_1 \ln \bar{p}$, where for the element ln f we define the right logarithmic derivative by $d_t(\ln f) := f^{-1} \dot{f}$. (cf. with [Mase et al., 2020])
- Then, the non-commutative derivatives are

$$\partial_q H = -\bar{p} + 2q^2 + t, \qquad \partial_{\bar{p}} H = -q - \bar{\alpha}_1 \, \bar{p}^{-1}, \qquad \partial_t H = q$$
(9)

and (8) is equivalent to the $d-P(E_7)$ system.

Affine Weyl groups and NC discrete systems

Affine Weyl groups W

- Let us fix a generalized Cartan matrix $C = (c_{ij})$, where $i, j \in I := \{0, 1, ..., n\}$.
- ► Sets $\Delta = \{\alpha_0, \dots, \alpha_n\}$, $\Delta^{\vee} = \{\alpha_0^{\vee}, \dots, \alpha_n^{\vee}\}$ correspond to simple roots and simple co-roots.
- ▶ Denote by Q = Q(C) and $Q^{\vee} = Q^{\vee}(C)$ the root and co-root lattices. The pairing $\langle \cdot, \cdot \rangle : Q \times Q^{\vee} \to \mathbb{Z}$ is defined by $\langle \alpha_i, \alpha_j^{\vee} \rangle = c_{ij}$ and $\alpha_i^{\vee} = 2\alpha_i / (\alpha_i, \alpha_i)$.
- ▶ Denote by W = W(C) the Weyl group (or the Coxeter group) defined by generators s_i , $i \in I$:

$$W(C) = \langle s_0, s_1, \dots, s_n \mid s_i^2 = 1, (s_i s_j)^{m_{ij}} = 1 \rangle,$$
(10)

where the exponents are determined by the value of the product $c_{ij}c_{ji}$ as below

$$\begin{array}{c|c} c_{ij}c_{ji} & 0 & 1 & 2 & 3 & \geq 4 \\ \hline m_{ij} & 2 & 3 & 4 & 6 & \infty \end{array}$$

These generators act naturally on Q by reflections

$$\mathbf{s}_i(\alpha_j) = \alpha_j - \langle \alpha_i, \alpha_j^{\vee} \rangle \, \alpha_i = \alpha_j - \mathbf{c}_{ij} \, \alpha_i. \tag{11}$$

- Each s_i-action on Q induces an automorphism of the field C(α) of rational functions in α_i. Hence, C(α) is a left W-module.
- Recall that one of the important properties of the affine Weyl groups is that they have translations, also known as Kac translations. Let W_0 be a finite Weyl group, $\delta = \sum_{i \in I} k_i \alpha_i$ be the null root and $V_0 = \{\mu \in V \mid \langle \mu, \delta^{\vee} \rangle = 0\}$. For an element $\mu \in V_0$ such that $\langle \mu, \mu^{\vee} \rangle \neq 0$ we define a translation element $t_{\mu} \in W$ by the formula

$$t_{\mu} = s_{\delta - \mu} \, s_{\mu} \tag{12}$$

and suppose that $w t_{\mu} = t_{w(\mu)} w$ for any $w \in W$.

Affine Weyl groups W

- Let us fix a generalized Cartan matrix $C = (c_{ij})$, where $i, j \in I := \{0, 1, \dots, n\}$.
- ► Sets $\Delta = \{\alpha_0, \dots, \alpha_n\}$, $\Delta^{\vee} = \{\alpha_0^{\vee}, \dots, \alpha_n^{\vee}\}$ correspond to simple roots and simple co-roots.
- ▶ Denote by Q = Q(C) and $Q^{\vee} = Q^{\vee}(C)$ the root and co-root lattices. The pairing $\langle \cdot, \cdot \rangle : Q \times Q^{\vee} \to \mathbb{Z}$ is defined by $\langle \alpha_i, \alpha_j^{\vee} \rangle = c_{ij}$ and $\alpha_i^{\vee} = 2\alpha_i/(\alpha_i, \alpha_i)$.
- ▶ Denote by W = W(C) the Weyl group (or the Coxeter group) defined by generators s_i , $i \in I$:

$$W(C) = \langle s_0, s_1, \dots, s_n \mid s_i^2 = 1, (s_i s_j)^{m_{ij}} = 1 \rangle,$$
(10)

where the exponents are determined by the value of the product $c_{ij}c_{ji}$ as below

These generators act naturally on Q by reflections

$$\mathbf{s}_i(\alpha_j) = \alpha_j - \langle \alpha_i, \alpha_j^{\vee} \rangle \, \alpha_i = \alpha_j - \mathbf{c}_{ij} \, \alpha_i. \tag{11}$$

- Each s_i-action on Q induces an automorphism of the field C(α) of rational functions in α_i. Hence, C(α) is a left W-module.
- ▶ Recall that one of the important properties of the affine Weyl groups is that they have translations, also known as Kac translations. Let W_0 be a finite Weyl group, $\delta = \sum_{i \in I} k_i \alpha_i$ be the null root and $V_0 = \{\mu \in V \mid \langle \mu, \delta^{\vee} \rangle = 0\}$. For an element $\mu \in V_0$ such that $\langle \mu, \mu^{\vee} \rangle \neq 0$ we define a translation element $t_{\mu} \in W$ by the formula

$$t_{\mu} = s_{\delta - \mu} \, s_{\mu} \tag{12}$$

and suppose that $w t_{\mu} = t_{w(\mu)} w$ for any $w \in W$.

Extended representations of W & NC discrete dynamics

The Kac translation acts on simple affine roots as follows

$$t_{\mu}(\alpha) = \alpha - \langle \mu, \alpha \rangle \,\delta = \alpha - \mu_{\alpha} \delta. \tag{13}$$

- It is known that the affine Weyl group is decomposed into a semi-direct product of translations in the lattice part *M* and the finite Weyl group *W*₀ acting on *M*, i.e. *W* = *M* × *W*₀. The lattice part *M* acts on C(α) as a shift operator, thanks to (13).
- Consider the set of elements $f_i \in \mathcal{R}$, $i \in I$, which we will often call functions or variables.
- ▶ We propose an extension of the representation of W on $\mathbb{C}(\alpha)$ to the skew field $\mathbb{C}(\alpha, f)$ of rational functions in α_i and elements f_i , $i \in I$. One needs to specify the action of s_i on f_j in such a way that the automorphisms s_i on $\mathbb{C}(\alpha, f)$ preserve the Weyl group structure.

Remark 1. Such classes of representations arise naturally from Bäcklund transformations of the differential Painlevé equations (e.g., see the review [Grammaticos and Ramani, 2004]).

For each $\mu \in M$ we define a set of elements $F_{\mu,i}(\alpha, f) \in \mathbb{C}(\alpha, f)$ by

$$t_{\mu}(f_i) = F_{\mu,i}(\alpha, f). \tag{14}$$

This set can be considered as a non-commutative discrete dynamical system.

Remark 2. Sometimes it is necessary to work with an extended Weyl group W. A similar description of the discrete dynamics can be given for \widetilde{W} as well.

Extended representations of W & NC discrete dynamics

The Kac translation acts on simple affine roots as follows

$$t_{\mu}(\alpha) = \alpha - \langle \mu, \alpha \rangle \, \delta = \alpha - \mu_{\alpha} \delta. \tag{13}$$

- It is known that the affine Weyl group is decomposed into a semi-direct product of translations in the lattice part *M* and the finite Weyl group *W*₀ acting on *M*, i.e. *W* = *M* ⋊ *W*₀. The lattice part *M* acts on C(α) as a shift operator, thanks to (13).
- Consider the set of elements $f_i \in \mathcal{R}$, $i \in I$, which we will often call functions or variables.
- We propose an extension of the representation of W on $\mathbb{C}(\alpha)$ to the skew field $\mathbb{C}(\alpha, f)$ of rational functions in α_i and elements f_i , $i \in I$. One needs to specify the action of s_i on f_j in such a way that the automorphisms s_i on $\mathbb{C}(\alpha, f)$ preserve the Weyl group structure.

Remark 1. Such classes of representations arise naturally from Bäcklund transformations of the differential Painlevé equations (e.g., see the review [Grammaticos and Ramani, 2004]).

For each $\mu \in M$ we define a set of elements $F_{\mu,i}(\alpha, f) \in \mathbb{C}(\alpha, f)$ by

$$t_{\mu}(f_i) = F_{\mu,i}(\alpha, f). \tag{14}$$

This set can be considered as a non-commutative discrete dynamical system.

Remark 2. Sometimes it is necessary to work with an extended Weyl group W. A similar description of the discrete dynamics can be given for \widetilde{W} as well.

Extended representations of *W* & NC discrete dynamics

▶ The Kac translation acts on simple affine roots as follows

$$t_{\mu}(\alpha) = \alpha - \langle \mu, \alpha \rangle \,\delta = \alpha - \mu_{\alpha} \delta. \tag{13}$$

- It is known that the affine Weyl group is decomposed into a semi-direct product of translations in the lattice part *M* and the finite Weyl group *W*₀ acting on *M*, i.e. *W* = *M* ⋊ *W*₀. The lattice part *M* acts on C(α) as a shift operator, thanks to (13).
- Consider the set of elements $f_i \in \mathcal{R}$, $i \in I$, which we will often call functions or variables.
- We propose an extension of the representation of W on $\mathbb{C}(\alpha)$ to the skew field $\mathbb{C}(\alpha, f)$ of rational functions in α_i and elements f_i , $i \in I$. One needs to specify the action of s_i on f_j in such a way that the automorphisms s_i on $\mathbb{C}(\alpha, f)$ preserve the Weyl group structure.

Remark 1. Such classes of representations arise naturally from Bäcklund transformations of the differential Painlevé equations (e.g., see the review [Grammaticos and Ramani, 2004]).

For each $\mu \in M$ we define a set of elements $F_{\mu,i}(\alpha, f) \in \mathbb{C}(\alpha, f)$ by

$$t_{\mu}(f_i) = F_{\mu,i}(\alpha, f). \tag{14}$$

This set can be considered as a non-commutative discrete dynamical system.

Remark 2. Sometimes it is necessary to work with an extended Weyl group \widetilde{W} . A similar description of the discrete dynamics can be given for \widetilde{W} as well.

NC d-Painlevé equations: from "continuous" to discrete

- Thanks to the paper [Bershtein et al., 2023], we know that matrix Hamiltonian Painlevé systems of all types have Bäcklund transformations forming an affine Weyl group structure.
- We have reconstructed all the generators for the extended affine Weyl groups corresponding to the NC Hamiltonian systems obtained in [Bobrova and Sokolov, 2023a].
- By using the translation operators similar to those presented in the papers [Sakai, 2001] or [Kajiwara et al., 2017], we have obtained a list of non-commutative discrete systems.
- ▶ They might be regarded as non-commutative analogs for the *d*-Painlevé systems.
- Note that they are connected by the degeneration procedure as follows



Remark. The d-P(E_6) and d-P(D_6) systems have a continuous limit to the P₂ system and its another non-equivalent version, respectively. The latter is a subcase of the system derived in [Adler and Sokolov, 2021] and labeled by P₂².

NC dressing chains and related discrete systems

NC dressing chains

Remark. The commutative dressing chain was introduced in [Veselov and Shabat, 1993]. It is related to the Painlevé equations [Adler, 1994] and arises from a generalisation of the symmetries for the P₄ and P₅ systems [Noumi and Yamada, 2000]. Quantum dressing chain might be found in [Nagoya, 2004]. Here we do not (!) assume any relations for the elements f_i .

▶ Let
$$j \in \mathbb{Z}/(n+1)\mathbb{Z}$$
. Consider the systems for $n = 2l$ and $n = 2l + 1$, $l \in \mathbb{Z}_{\geq 0}$, respectively

$$\begin{split} \dot{f}_{j} &= \sum_{1 \leq r \leq l} f_{j} f_{j+2r-1} - \sum_{1 \leq r \leq l} f_{j+2r} f_{j} + \alpha_{j}; \\ \frac{1}{2} t \dot{f}_{j} &= \sum_{1 \leq r \leq s \leq l} f_{j} f_{j+2r-1} f_{j+2s} - \sum_{1 \leq r \leq s \leq l} f_{j+2r} f_{j+2s+1} f_{j} \\ &+ \left(\frac{1}{2} - \sum_{1 \leq r \leq l} \alpha_{j+2r}\right) f_{j} + \alpha_{j} \sum_{1 \leq r \leq l} f_{j+2r}. \end{split}$$

• We will cal them $A_n^{(1)}$, $n \ge 2$ type systems or dressing chains in the Noumi-Yamada variables.

These systems admit Lax pairs.

Lax pairs

• Let $\Psi = \Psi(\lambda, t) \in Mat_{n+1}(\mathbb{R})$, $\lambda \in \mathbb{Z}(\mathbb{R})$ satisfy the linear system

$$\begin{cases} \partial_{\lambda}\Psi(\lambda, t) = \mathcal{A}(\lambda, t)\Psi(\lambda, t), \\ \partial_{t}\Psi(\lambda, t) = \mathcal{B}(\lambda, t)\Psi(\lambda, t), \end{cases}$$
(15)

where matrices $\mathcal{A} = \mathcal{A}(\lambda, t)$ and $\mathcal{B} = \mathcal{B}(\lambda, t)$ belong to $Mat_{n+1}(\mathcal{R})$ and have the form

$$\mathcal{A}(\lambda) = A_0 + A_{-1} \lambda^{-1}, \qquad \qquad \mathcal{B}(\lambda) = B_1 \lambda + B_0. \tag{16}$$

• Consider the matrices expressed in terms of the standard unit matrices $E_{r,s} \in Mat_{n+1}(\mathbb{C})$ as

$$A_{0} = E_{1,n} + f_{0} E_{1,n+1} + E_{2,n+1}, \quad A_{-1} = \sum_{1 \le r \le n+1} \beta_{r} E_{r,r} + \sum_{1 \le r \le n} f_{r} E_{r+1,r} + \sum_{1 \le r \le n-1} E_{r+2,r},$$

$$B_1 = E_{1,n+1}, \quad B_0 = \sum_{1 \le r \le n+1} g_r E_{r,r} + \sum_{1 \le r \le n} E_{r+1,r}.$$

• Let $\alpha_0 = 1 + \beta_{n+1} - \beta_1$, $\alpha_j = \beta_j - \beta_{j+1}$, $j \in \mathbb{Z}/(n+1)\mathbb{Z} \setminus \{0\}$.

Theorem. [Bobrova, 2024] There exists a set of the *g*-functions such that the compatibility condition of system (15) is equivalent to either the $A_{2l}^{(1)}$ or $A_{2l+1}^{(1)}$ system.

For the $A_{2i}^{(1)}$, we have $g_j = -\sum_{1 \le r \le l} f_{j+2r}$, where indexes belong to $\mathbb{Z}/(n+1)\mathbb{Z}$.

Lax pairs

• Let $\Psi = \Psi(\lambda, t) \in Mat_{n+1}(\mathbb{R})$, $\lambda \in \mathbb{Z}(\mathbb{R})$ satisfy the linear system

$$\begin{cases} \partial_{\lambda}\Psi(\lambda, t) = \mathcal{A}(\lambda, t)\Psi(\lambda, t), \\ \partial_{t}\Psi(\lambda, t) = \mathcal{B}(\lambda, t)\Psi(\lambda, t), \end{cases}$$
(15)

where matrices $\mathcal{A} = \mathcal{A}(\lambda, t)$ and $\mathcal{B} = \mathcal{B}(\lambda, t)$ belong to $Mat_{n+1}(\mathcal{R})$ and have the form

$$\mathcal{A}(\lambda) = A_0 + A_{-1} \lambda^{-1}, \qquad \qquad \mathcal{B}(\lambda) = B_1 \lambda + B_0. \tag{16}$$

• Consider the matrices expressed in terms of the standard unit matrices $E_{r,s} \in Mat_{n+1}(\mathbb{C})$ as

$$A_{0} = E_{1,n} + f_{0} E_{1,n+1} + E_{2,n+1}, \quad A_{-1} = \sum_{1 \le r \le n+1} \beta_{r} E_{r,r} + \sum_{1 \le r \le n} f_{r} E_{r+1,r} + \sum_{1 \le r \le n-1} E_{r+2,r},$$

$$B_1 = E_{1,n+1}, \quad B_0 = \sum_{1 \le r \le n+1} \frac{g_r}{g_r} E_{r,r} + \sum_{1 \le r \le n} E_{r+1,r}$$

• Let $\alpha_0 = 1 + \beta_{n+1} - \beta_1$, $\alpha_j = \beta_j - \beta_{j+1}$, $j \in \mathbb{Z}/(n+1)\mathbb{Z} \setminus \{0\}$.

Theorem. [Bobrova, 2024] There exists a set of the *g*-functions such that the compatibility condition of system (15) is equivalent to either the $A_{2l}^{(1)}$ or $A_{2l+1}^{(1)}$ system.

For the $A_{2l}^{(1)}$, we have $g_j = -\sum_{1 \le r \le l} f_{j+2r}$, where indexes belong to $\mathbb{Z}/(n+1)\mathbb{Z}$.

Bäcklund transformations and NC discrete dynamics

- Let the Cartan matrix C be of type $A_n^{(1)}$, $n \ge 2$ and $I = \{0, 1, \dots, n\}$.
- Let us set

$$\begin{aligned} s_i(\alpha_i) &= -\alpha_i, \qquad s_i(\alpha_j) = \alpha_j + \alpha_i \qquad (j = i \pm 1), \qquad s_i(\alpha_j) = \alpha_j \quad (j \neq i \pm 1), \\ s_i(f_i) &= f_i, \qquad s_i(f_j) = f_j \pm \alpha_i f_i^{-1} \qquad (j = i \pm 1), \qquad s_i(f_j) = f_j \qquad (j \neq i \pm 1), \\ \pi(\alpha_j) &= \alpha_{j+1}, \qquad \pi(f_j) = f_{j+1}, \qquad j \in \mathbb{Z}/(n+1)\mathbb{Z}. \end{aligned}$$

Theorem. [Bobrova, 2024] Transformations given above are Bäcklund transformations of the $A_{2l}^{(1)}$ and $A_{2l+1}^{(1)}$ systems. Moreover, they define a birational representation of the extended affine Weyl group of type $A_n^{(1)}$, $n \ge 2$.

Note that the shift operators are given by

 $T_1 = \pi s_n s_{n-1} \dots s_1, \quad T_2 = s_1 \pi s_n \dots s_2, \quad \dots, \quad T_{n+1} = s_n \dots s_1 \pi.$ (17)

- They satisfy the relation $T_1 T_2 \ldots T_{n+1} = 1$.
- Thus, any n of them form a basis for the lattice and we can define a discrete system.

Remark. Cases n = 2 and n = 3 correspond to the P₄ and P₅ equations and discrete systems labeled by d-P(E_6) and d-P(D_5) respectively. For n = 1 one needs to consider the P₂ system.

Bäcklund transformations and NC discrete dynamics

- Let the Cartan matrix C be of type $A_n^{(1)}$, $n \ge 2$ and $I = \{0, 1, \dots, n\}$.
- Let us set

$$\begin{aligned} s_i(\alpha_i) &= -\alpha_i, \qquad s_i(\alpha_j) = \alpha_j + \alpha_i \qquad (j = i \pm 1), \qquad s_i(\alpha_j) = \alpha_j \quad (j \neq i \pm 1), \\ s_i(f_i) &= f_i, \qquad s_i(f_j) = f_j \pm \alpha_i f_i^{-1} \qquad (j = i \pm 1), \qquad s_i(f_j) = f_j \qquad (j \neq i \pm 1), \\ \pi(\alpha_j) &= \alpha_{j+1}, \qquad \pi(f_j) = f_{j+1}, \qquad j \in \mathbb{Z}/(n+1)\mathbb{Z}. \end{aligned}$$

Theorem. [Bobrova, 2024] Transformations given above are Bäcklund transformations of the $A_{2l}^{(1)}$ and $A_{2l+1}^{(1)}$ systems. Moreover, they define a birational representation of the extended affine Weyl group of type $A_n^{(1)}$, $n \ge 2$.

Note that the shift operators are given by

 $T_1 = \pi \, s_n \, s_{n-1} \, \dots \, s_1, \quad T_2 = s_1 \, \pi \, s_n \, \dots \, s_2, \quad \dots, \quad T_{n+1} = s_n \, \dots \, s_1 \, \pi. \tag{17}$

- They satisfy the relation $T_1 T_2 \ldots T_{n+1} = 1$.
- Thus, any n of them form a basis for the lattice and we can define a discrete system.

Remark. Cases n = 2 and n = 3 correspond to the P₄ and P₅ equations and discrete systems labeled by d-P(E_6) and d-P(D_5) respectively. For n = 1 one needs to consider the P₂ system.

Further questions

Further questions

A study of the NC *d*-Painlevé equations

- ▶ We expect that these equations admit Lax pairs and have a Hamiltonian form.
- We also expect that they have a continuous limit to known NC "differential" Painlevé systems obtained in [Bobrova and Sokolov, 2023b].
- Commutative d- and differential Painlevé equations are related to the orthogonal polynomials [Van Assche, 2022]. Orthogonal polynomials have a non-commutative analog (see [Gelfand et al., 1995]). We assume that our equations are connected with them.

Other discrete Painlevé equations

- Our method can be applied to the *q*-discrete Painlevé equations.
- In particular, one may define a NC version of the q-P₆ equation which generalizes the matrix equation obtained in [Kawakami, 2020] to the purely non-commutative case. (an ongoing project)
- What about a non-commutative *ell*-discrete Painlevé equation? (see also [Okounkov and Rains, 2015])

NC geometry related to Painlevé equations

- ▶ What is the Okamoto space of initial data of NC "differential" Painlevé equations?
- ▶ We would like to generalize the method of Painlevé equations' classification introduced in Sakai's paper [Sakai, 2001]. Recent developments might be found in [Rains, 2019].

Cluster algebras and discrete Painlevé equations

It is known that discrete Painlevé equations are connected with cluster algebras (see, e.g. [Bershtein et al., 2018]). Might we have the same connection in a NC case?

Many thanks!

References I

[Adler, 1994] Adler, V. E. (1994). Nonlinear chains and Painlevé equations. Physica D: Nonlinear Phenomena, 73(4):335–351.

[Adler, 2020] Adler, V. E. (2020). Painlevé type reductions for the non-Abelian Volterra lattices. Journal of Physics A: Mathematical and Theoretical, 54(3):035204.

arXiv:2010.09021.

[Adler and Sokolov, 2021] Adler, V. E. and Sokolov, V. V. (2021). On matrix Painlevé II equations. Theoret. and Math. Phys., 207(2):188–201. arXiv:2012.05639.

[Bershtein et al., 2018] Bershtein, M., Gavrylenko, P., and Marshakov, A. (2018). Cluster integrable systems, q-Painlevé equations and their quantization. Journal of High Energy Physics, 2018(2):1 – 29. arXiv:1711.02063.

[Bershtein et al., 2023] Bershtein, M., Grigorev, A., and Shchechkin, A. (2023). Hamiltonian reductions in matrix Painlevé systems. Letters in Mathematical Physics, 113(2):47. arXiv:2208.04824.

[Bertola et al., 2018] Bertola, M., Cafasso, M., and Rubtsov, V. (2018). Noncommutative Painlevé equations and systems of Calogero type. Communications in Mathematical Physics, 363(2):503–530. arXiv:1710.00736.

References II

[Bobrova, 2024] Bobrova, I. (2024).

Affine Weyl groups and non-Abelian discrete systems: an application to the *d*-Painlevé equations. *Journal of Nonlinear Science (under review)*. arXiv:2403.18463.

[Bobrova and Sokolov, 2023a] Bobrova, I. and Sokolov, V. (2023a). Classification of Hamiltonian non-abelian Painlevé type systems. *Journal of Nonlinear Mathematical Physics*, 30:646–662. arXiv:2209.00258.

[Bobrova and Sokolov, 2023b] Bobrova, I. and Sokolov, V. (2023b). On classification of non-abelian Painlevé type systems. Journal of Geometry and Physics, 191:104885. arXiv:2303.10347.

[Cafasso and Manuel, 2014] Cafasso, M. and Manuel, D. (2014). Non-commutative Painlevé equations and Hermite-type matrix orthogonal polynomials. Communications in Mathematical Physics, 326(2):559–583. arXiv:1301.2116.

[Cafasso et al., 2018] Cafasso, M., Manuel, D., et al. (2018).

The Toda and Painlevé systems associated with semiclassical matrix-valued orthogonal polynomials of Laguerre type. *SIGMA. Symmetry, Integrability and Geometry: Methods and Applications*, 14:076. arXiv:1801.08740.

References III

[Cassatella-Contra et al., 2014] Cassatella-Contra, G. A., Manas, M., and Tempesta, P. (2014). Singularity confinement for matrix discrete Painlevé equations. Nonlinearity, 27(9):2321.

[Cohn, 1995] Cohn, P. M. (1995).

Skew Fields: Theory of General Division Rings.

Encyclopedia of Mathematics and its Applications. Cambridge University Press.

[Gambier, 1910] Gambier, B. (1910).

Sur les équations différentielles du second ordre et du premier degré dont l'intégrale générale est à points critiques fixes.

Acta Mathematica, 33(1):1-55.

[Gelfand et al., 1995] Gelfand, I. M., Krob, D., Lascoux, A., Leclerc, B., Retakh, V. S., and Thibon, J.-Y. (1995). Noncommutative Symmetric Functions. Advances in Mathematics, 2(112):218–348. arXiv:hep-th/9407124.

[Grammaticos and Ramani, 2004] Grammaticos, B. and Ramani, A. (2004). Discrete painlevé equations: a review. Discrete integrable systems, pages 245–321.

[Kajiwara et al., 2017] Kajiwara, K., Noumi, M., and Yamada, Y. (2017). Geometric aspects of Painlevé equations. Journal of Physics A: Mathematical and Theoretical, 50(7):073001. arXiv:1509.08186.

References IV

[Kawakami, 2015] Kawakami, H. (2015). Matrix Painlevé systems. Journal of Mathematical Physics, 56(3):033503.

[Kawakami, 2020] Kawakami, H. (2020).
 A *q*-analogue of the matrix sixth Painlevé system.
 Journal of Physics A: Mathematical and Theoretical, 53(49):495203.
 arXiv:2301.12837.

[Kontsevich, 1993] Kontsevich, M. (1993). Formal (non)-commutative symplectic geometry, The Gelfand Mathematical Seminars, 1990–1992. Fields Institute Communications. Birkhäuser Boston, pages 173–187.

[Mase et al., 2020] Mase, T., Nakamura, A., and Sakai, H. (2020).
 Discrete Hamiltonians of discrete Painlevé equations.
 In Annales de la Faculté des sciences de Toulouse: Mathématiques, volume 29 (5), pages 1251–1264.

[Nagoya, 2004] Nagoya, H. (2004).

Quantum Painlevé systems of type $A_{l}^{(1)}$. International Journal of Mathematics, 15(10):1007–1031. arXiv:math/0402281v2.

[Noumi and Yamada, 1998] Noumi, M. and Yamada, Y. (1998). Affine Weyl groups, discrete dynamical systems and Painlevé equations. Communications in Mathematical Physics, 199:281–295.

References V

[Noumi and Yamada, 2000] Noumi, M. and Yamada, Y. (2000). Affine Weyl group symmetries in Painlevé type equations. *Citeseer.*

[Okounkov and Rains, 2015] Okounkov, A. and Rains, E. (2015). Noncommutative geometry and Painlevé equations. Algebra Number Theory, 9(6):1363–1400. arXiv:1404.5938.

[Olver and Sokolov, 1998] Olver, P. J. and Sokolov, V. V. (1998). Integrable evolution equations on associative algebras. Communications in Mathematical Physics, 193(2):245–268.

[Painlevé, 1902] Painlevé, P. (1902).

Sur les équations différentielles du second ordre et d'ordre supérieur dont l'intégrale générale est uniforme. Acta mathematica, 25:1–85.

[Rains, 2019] Rains, E. M. (2019).

The birational geometry of noncommutative surfaces. arXiv preprint arXiv:1907.11301.

[Retakh and Rubtsov, 2010] Retakh, V. S. and Rubtsov, V. N. (2010). Noncommutative Toda Chains, Hankel Quasideterminants and Painlevé II Equation. Journal of Physics A, Mathematical and Theoretical, 43(50):505204. arXiv:1007.4168.

References VI

[Sakai, 2001] Sakai, H. (2001).

Rational surfaces associated with affine root systems and geometry of the Painlevé equations. *Communications in Mathematical Physics*, 220(1):165–229.

[Van Assche, 2022] Van Assche, W. (2022). Orthogonal polynomials, Toda lattices and Painlevé equations. *Physica D: Nonlinear Phenomena*, 434:133214. arXiv:2202.11017.

[Veselov, 1991] Veselov, A. P. (1991).
 Integrable maps.
 Russian Mathematical Surveys, 46(5):1.

[Veselov and Shabat, 1993] Veselov, A. P. and Shabat, A. B. (1993). Dressing chains and the spectral theory of the Schrödinger operator. *Functional Analysis and Its Applications*, 27(2):81–96.